# Random positive maps 

Benoît Collins

Kyoto University \& University of Ottawa
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Joint work in preparation with Patrick Hayden (Stanford) and Ion Nechita (CNRS, Toulouse)

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1. Positive maps: why do we care? (a primer of quantum information theory)
2. Random positive maps with random matrices (convergence of the largest eigenvalue).
3. Application: almost optimal entanglement witnesses.

## Quantum information: a primer

- A quantum system is a Hilbert space $\mathbb{C}^{n}$. Its set of states $D\left(\mathbb{C}^{n}\right)=D_{n}$ is the collection of positive trace one matrices of $\mathcal{M}_{n}(\mathbb{C})$.


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- Entangled states $\operatorname{Ent}\left(n_{1}, n_{2}\right):=D\left(n_{1} n_{2}\right)-\operatorname{Sep}\left(n_{1}, n_{2}\right)$. A very important set (resource for quantum computing, etc).


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- $\Phi_{1}$ is $k$-positive for all $k$ whereas $\Phi_{2}$ is 'only' 1-positive.
- A map that is positive for all $k$ is called completely positive.


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- On the other hand, positive maps are still completely unclassified.
(roughly speaking) The only final results available are: maps from $\mathcal{M}_{n_{1}}(\mathbb{C}) \rightarrow \mathcal{M}_{n_{2}}(\mathbb{C})$ with $\left(n_{1}, n_{2}\right)=\{(1, n) ;(n, 1) ;(2,2) ;(2,3) ;(3,2)\}$ are positive iff they are CP.


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- However, if $\rho \in \operatorname{Ent}\left(n_{1}, n_{2}\right)$ then $\Phi \otimes I_{n_{2}}(\rho)$ could in principle fail to be positive.
- But a failure to be positive can't happen if $\Phi$ is CP by definition.
So, trying to find positive but not completely positive maps is a strategy to witness entanglement.


## Quantum information: entanglement witness

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- The following state is not separable in $\mathcal{M}_{2}(\mathbb{C}) \otimes \mathcal{M}_{2}(\mathbb{C})$

$$
\left(\begin{array}{cccc}
0.2 & 0 & 0 & 0 \\
0 & 0.3 & 0.3 & 0 \\
0 & 0.3 & 0.3 & 0 \\
0 & 0 & 0 & 0.2
\end{array}\right)
$$

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- RMT can help.


## Choi matrix

- To a linear map $\Phi: \mathcal{M}_{n_{1}}(\mathbb{C}) \rightarrow \mathcal{M}_{n_{2}}(\mathbb{C})$ we assocciate its Choi matrix $C_{\Phi} \in \mathcal{M}_{n_{1}}(\mathbb{C}) \otimes \mathcal{M}_{n_{2}}(\mathbb{C})$ given by

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- More recently: $\Phi$ is positive iff $p \otimes 1_{n_{2}} C_{\Phi} p \otimes 1_{n_{2}}$ is positive.


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And if $a$ is such that $1-2 \sqrt{a / k}>0$, with probability tending to 1 as $n$ becomes large, we obtain a random positive map.
[largest eigenvalue convergence $+\varepsilon$-net + union bound argument]
- In addition, if $1-2 \sqrt{a}<0, \Phi$ is not completely positive with probability tending to 1 as $n$ becomes large, therefore it 'detects' many entangled states.


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Let $X$ be a $G U E$ as above: $X \in \mathcal{M}_{k}(\mathbb{C}) \otimes \mathcal{M}_{n}(\mathbb{C})$ its eigenvalues are approximately in $[1-2 \sqrt{a}, 1+2 \sqrt{a}] . k$ is fixed, $n$ tends to $\infty$.

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- $X$ is positive with probability one as $n \rightarrow \infty$ as soon as $0<\alpha<1$.
- $X$ is PPT with probability one as $n \rightarrow \infty$ as soon as $0<\alpha<1$.
- PPT states and general states have typical size - i.e. PPT is not so efficient in large dimension to detect entanglement.


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- In both cases, $\alpha$ is of order $C / \sqrt{k}$. The order $C / \sqrt{k}$ is optimal. We use the non-centered GUE random positive maps exhibited earlier to prove this result.


## How useful are these examples?

Let $X$ be a $G U E$ as above: $X \in \mathcal{M}_{k}(\mathbb{C}) \otimes \mathcal{M}_{n}(\mathbb{C})$ its eigenvalues are approximately in $[1-2 \sqrt{a}, 1+2 \sqrt{a}] . k$ is fixed, $n$ tends to $\infty$. Set $\alpha=2 \sqrt{a}$.
With probability one, as $n \rightarrow \infty$ :

- With $1>\alpha>4 / \sqrt{k}, X$ is PPT but not separable.
- If $\alpha<\sqrt{k} /(2(k-1)+\sqrt{k}), X$ is separable.
- The criterion starts to become useful when $k>16$.
- In both cases, $\alpha$ is of order $C / \sqrt{k}$. The order $C / \sqrt{k}$ is optimal. We use the non-centered GUE random positive maps exhibited earlier to prove this result.
- Conclusion: Random maps are much more efficient than PPT.

Thank you!

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